

I) 1) $\vec{e}_\rho = -\cos\psi \vec{e}_x + \sin\psi \vec{e}_y$

Une force est conservative si son travail est independant du chemin suivi.

1.3) $E_c = \frac{1}{2} m v^2 = \frac{1}{2} m \dot{\rho}^2 + \frac{1}{2} m \rho^2 \dot{\psi}^2$
 $E_m = E_c + E_p$

1.4) $T_{PM} = \frac{dE_m}{dt}$

1.5) $\vec{P} = m\vec{g} = -mg\vec{e}_z$

1.6) $dL(H) = R d\psi \vec{e}_\rho = \frac{dL}{dt} = R \dot{\psi} \vec{e}_\rho$

1.7) $E_c = \frac{1}{2} m \dot{\rho}^2 + m g R (1 - \cos\psi)$

1.8) Seule la reaction est non conservative. Mais $R_f \perp$ déplacement \Rightarrow son travail est nul $\Rightarrow \mathcal{L} = 0 \Rightarrow \frac{dE}{dt} = 0$: E. conserve.

1.9) $T_{PM} = m R \dot{\psi}^2 + m g R \sin\psi \cdot \dot{\psi} = 0 \Leftrightarrow \dot{\psi} + \frac{g}{R} \sin\psi = 0$

1.10) $\psi(t=0) = 0 \Rightarrow A = 0 \Rightarrow \dot{\psi}(t) = B \cos\omega t$

1.11) $v(t=0) = v_0 = R \dot{\psi}(t=0) = B R \omega \Rightarrow B = \frac{v_0}{R \omega}$

1.12) $T_{EC} = \Delta E_c = W(\text{travaux}) = \int \vec{P} \cdot d\vec{l} = -mg(R - R \cos\psi) = -mg(R + z)$

1.13) $v = 0 \Leftrightarrow 2g(R + z) = v_0^2 \Leftrightarrow z_0 = \frac{v_0^2}{2g} - R$

Si H arrive en z = z_0, alors il peut faire un tour complet.
 $\Rightarrow R = \frac{v_0^2}{2g} - R \Leftrightarrow (v_0)_c = 2\sqrt{Rg}$

I) C-1) $\vec{\omega}(P, A) = \omega \vec{e}_y$

1.2) $\vec{v}(M, P) = \left[\frac{dOM}{dt} \right] \vec{e}_x = R \dot{\psi} \vec{e}_y$

1.3) $\vec{v}_E(M, P) = \vec{v}(O, E, P) + \vec{\omega}(P, A) \times \vec{OM} = \omega R \vec{e}_y \times R \vec{e}_\rho = \omega R \sin(\pi - \psi) \vec{e}_x$

1.4) $\vec{v}(M, P) = \vec{v}(M, P) + \vec{v}(H, P) = \omega R \vec{e}_y + \omega R \sin\psi \vec{e}_x$

1.5) $\vec{a}(M, P) = R(\ddot{\psi} \vec{e}_y - \dot{\psi}^2 \vec{e}_\rho)$

1.6) $\vec{a}(M, P) = 2\omega \vec{e}_y + R \ddot{\psi} \vec{e}_y - 2\omega \dot{\psi} \cos\psi \vec{e}_x$

1.7) $\vec{a}(M, P) = \vec{a}(M, P) + \vec{a}(H, P) = \frac{R}{\rho}(\ddot{\psi} \vec{e}_y - \dot{\psi}^2 \vec{e}_\rho) + \omega R \sin\psi \vec{e}_x$

1.8) $\vec{P} + \vec{R}_f = m \vec{a} \Rightarrow$ on projette sur \vec{e}_ρ
 $\Rightarrow R_f = -mg \cos\psi + m[-R\dot{\psi}^2 + \omega^2 R \sin^2\psi]$

II) 1) $\lambda_p = \frac{0.15}{1-p}$
 $\lambda_n = \frac{0.15}{1-p} = \frac{0.15}{1-p}$

2) Si le cas d'une force centrale, l'axe balayé par le rayon vecteur pendant un intervalle de temps donné est etc.

2.3) $\frac{dS}{dt} = \frac{1}{2} \|\vec{r} \times \vec{v}\| \Rightarrow C = r^2 \dot{\theta} = r^2 \dot{\theta}$

2.4) $\frac{dC}{dt} = 2r\dot{r}\dot{\theta} + r^2\ddot{\theta} = 0 \Rightarrow \ddot{\theta} = -\frac{2\dot{r}\dot{\theta}}{r^2}$

2.5) $\vec{a} = \ddot{r} \vec{e}_r + r \ddot{\theta} \vec{e}_\theta \Rightarrow \vec{a} = (\ddot{r} - r \dot{\theta}^2) \vec{e}_r + (2\dot{r}\dot{\theta} + r \ddot{\theta}) \vec{e}_\theta$

2.6) $\ddot{r} = \frac{r}{1-\epsilon \cos\theta} \Rightarrow \ddot{r} = \frac{r(1-\epsilon)(-\sin\theta)\dot{\theta}}{(1-\epsilon \cos\theta)^2} = -\frac{r \epsilon \sin\theta}{(1-\epsilon \cos\theta)^2} \dot{\theta}$

2.7) $\ddot{r} = -\frac{r \epsilon \sin\theta}{(1-\epsilon \cos\theta)^2} \dot{\theta} \Rightarrow \ddot{r} = -\frac{r \epsilon \sin\theta}{r^2} \dot{\theta} = -\frac{\epsilon \sin\theta}{r} \dot{\theta}$

2.8) $\vec{a} = -\frac{r \epsilon \cos\theta}{r^2} \dot{\theta} = -\frac{\epsilon \cos\theta}{r} \dot{\theta} = -\left(\frac{\epsilon \cos\theta}{r} + \frac{r \dot{\theta}^2}{r} \right) \vec{e}_r$

1.9) $\vec{F} = m \vec{a} = -m \frac{\epsilon \cos\theta}{r} \dot{\theta} \vec{e}_r =$ force centrale - c'est la force d'interaction gravitationnelle.

